

EFT & GR #2

Note Title

Review of end of lecture #1:

$$\mathcal{L}_{eff} = \frac{v^2}{4} \text{Tr} (\partial_\mu U \partial^\mu U^\dagger) + l_1 [\text{Tr} (\partial_\mu U \partial^\mu U^\dagger)]^2 + l_2 \text{Tr} (\partial_\mu U \partial_\nu U^\dagger) \text{Tr} (\partial^\mu U \partial^\nu U^\dagger)$$

9/10/13

$$U = \exp\left(i \frac{\vec{L} \cdot \vec{\Pi}}{v}\right)$$

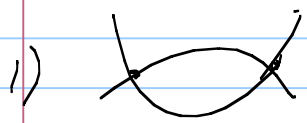
$$\begin{aligned} \mathcal{M}_{full} = & \frac{t}{v^2} + \left[\frac{1}{m_\sigma^2 v^2} - \frac{11}{96\pi^2 v^4} \right] t^2 \\ & - \frac{1}{144\pi^2 v^4} [s(s-u) + u(u-s)] \\ & - \frac{1}{96\pi^2 v^4} \left[3t^2 \ln \frac{-t}{m_\sigma^2} + s(s-u) \ln \frac{-s}{m_\sigma^2} + u(u-s) \ln \frac{-u}{m_\sigma^2} \right] \end{aligned} \quad (3.7)$$

ie effective theory result [Le72, GaL84] has a very similar form but does not know about the existence of the σ ,

$$\begin{aligned} \mathcal{M}_{eff} = & \frac{t}{v^2} + \left[8\ell_1 + 2\ell_2 + \frac{5}{192\pi^2} \right] \frac{t^2}{v^4} \\ & + \left[2\ell_2 + \frac{7}{576\pi^2} \right] [s(s-u) + u(u-s)]/v^4 \\ & - \frac{1}{96\pi^2 v^4} \left[3t^2 \ln \frac{-t}{\mu^2} + s(s-u) \ln \frac{-s}{\mu^2} + u(u-s) \ln \frac{-u}{\mu^2} \right] \end{aligned} \quad (3.8)$$

Matching

Comments:



$$l_1^{(2)} = l_1 + \frac{1}{384\pi^2} \left[\frac{2}{4-d} + \dots \right]$$

renormalized
 E^4

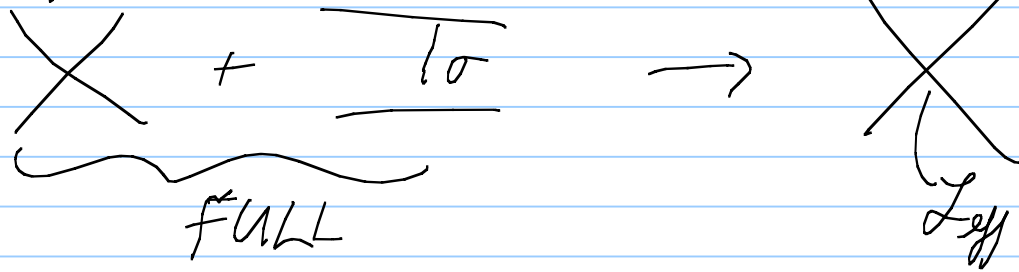
2) Structure same

3) Analytic term s^2, t^2, \dots comes with $h_i \leftarrow \frac{10^2}{8M_p^2}$

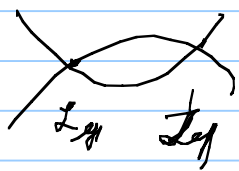
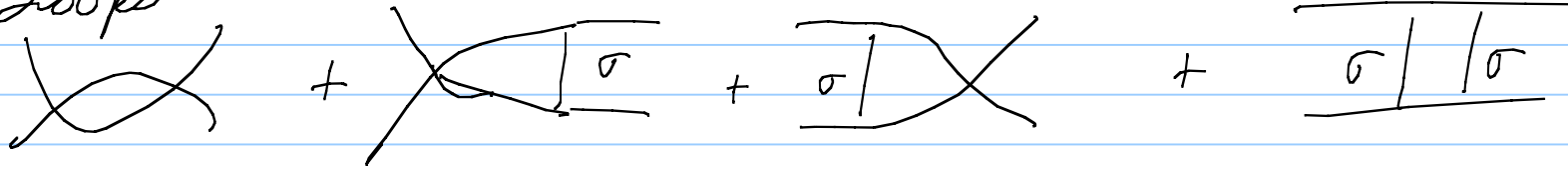
4) Nonanalytic $\ln t$ - no coeff. long distances

5) Loops E^4

Why does it work?



Loops



→ Low loop momentums - same by construction

→ High loop momenta - goes local l₁

Power counting

$$M \sim \left(\frac{t}{N^2} \right) \sim \mathcal{O}(E^2)$$

$$\text{Loops} \sim \mathcal{O}(E^4) \quad \mathcal{L}_4 \sim \mathcal{O}(E^4)$$

$$I_{\text{loop}} = \int \frac{d^4 l}{(2\pi)^4} \frac{E^2}{N^2} \quad \text{Propagators} \quad \frac{E^2}{N^2}$$

powers of N \rightarrow $\frac{E^4}{N^4}$ $\xrightarrow{\text{Regularize dimensionally}}$

Weinberg's power counting theorem

$$E^2 \rightarrow \text{Tree level}$$


$$E^4 \rightarrow L_4 + 1 \text{ loop}$$

$$E^6 \rightarrow L_6 + 2 \text{ loop } L_2 + 1 \text{ loop } (L_2 + L_4)$$

Limits of EFT

$$M = \frac{t}{\Lambda^2} + \frac{t^2}{\Lambda^2} + \frac{t^2}{\Lambda^2} \ln t$$

light σ \swarrow \nwarrow loop



"Rules:"

- Most general \mathcal{L} consistent with symmetries
- Order by energy expansion
- Start with lowest order \mathcal{L}
- Loop
- renormalize parameters \rightarrow match to full theory
 \rightarrow measure
- predictions

Heat kernel

$$H(X, \tilde{\tau}) = \langle X | e^{-\tilde{\tau} D} | X \rangle$$

Parameter

$$D = d_\mu d^\mu + m^2 + \sigma^2$$

$$d_\mu = \partial_\mu + \Gamma_\mu$$

$$H(X, \tilde{\tau}) = \frac{i}{(4\pi)^{d/2}} \frac{e^{-\tilde{\tau} m^2}}{\tilde{\tau}^{d/2}} \left[a_0(X) + a_1(X) \tilde{\tau} + a_2(X) \tilde{\tau}^2 + \dots \right]$$

$$\det D = e^{\text{tr} \ln D} = e^{\int d^d X \text{Tr} \langle X | \ln D | X \rangle}$$

$$\ln \frac{b}{a} = \int_0^\infty \frac{dx}{x} (e^{-ax} - e^{-bx})$$

$$\langle X | \ln D | X \rangle = -\frac{i}{(4\pi)^{d/2}} \sum_{m=0}^{\infty} m^{d-2m} a_m(x) \Gamma(m - d/2)$$

$\leftarrow a_2 \sim \Gamma(2 - d/2)$

$$H(x, \tilde{z}) = \int \frac{d^d p}{(2\pi)^d} e^{-ip \cdot X} e^{-\tilde{z} D} e^{ipx}$$

$$\langle X | P \rangle = e^{ipx}$$

$$d_m e^{ipx} = e^{ipx} (ip_m + d_m)$$

$$\therefore a_0(x) = 1, \quad a_1(x) = -\sigma$$

$$a_2(x) = \frac{1}{2} \sigma^2 + \frac{1}{12} [d_m, d_\nu] [d^m, d^\nu] + \frac{1}{6} [d_m, [\bar{d}^m, \sigma]]$$



Background field

$$U = \bar{U} + \delta U$$

$$U^\dagger U = 1$$

$$U = \bar{U} e^{i\Delta}$$

$$\Delta = \tau^i \Delta^i$$

↙ E_0 of \mathcal{M}

$$\begin{aligned} \text{Tr}(\partial_\mu U \partial^\mu U^\dagger) &= \text{Tr}(\partial_\mu \bar{U} \partial^\mu \bar{U}^\dagger) - 2i \text{Tr}(\bar{U}^\dagger \partial_\mu \bar{U} \partial^\mu \Delta) \\ &\quad + \text{Tr}(\partial^\mu \Delta \partial_\mu \Delta + \bar{U} \partial^\mu \bar{U} (\Delta \partial^\mu \Delta - (\partial_\mu \Delta) \Delta)) \end{aligned}$$

$$\mathcal{L} = \mathcal{L}(\bar{U}) - \frac{N^2}{2} \Delta^i (\partial_\mu \Delta^\mu + \sigma) \tau^{ij} \Delta^j$$

$$\tau^{ij} = -\frac{1}{4} \text{Tr}([\tau^i, \tau^j] \bar{U}^\dagger \partial_\mu U), \quad \sigma = \text{Tr}([\tau^i, \bar{U}^\dagger \partial_\mu \bar{U}] [\tau^j, \bar{U}^\dagger \partial^\mu \bar{U}])$$

+ Heat Kernel

$$l_1^R = l_1 + \frac{1}{384\pi^2} \left[\frac{2}{d-4} + \dots \right]$$

$$\delta S^{1\text{loop}} = \frac{1}{32\pi^2} \left[\frac{2}{d-4} + \dots \right] \int d^4x \left[\frac{1}{12} \left[\text{Tr}(\partial_\mu \bar{u}^+ \partial^\mu \bar{u}) \right]^2 + \frac{1}{6} \text{Tr}(\partial_\mu \bar{u}^+ \partial_\nu \bar{u}) \right. \\ \left. + \pi \times \text{Tr}(\partial_\mu \bar{u}^+ \partial_\nu \bar{u}) \right]$$

\Rightarrow get divergences + renormalize

QCD full theory ψ, G , eff FT (π, K, η)

$$\psi_{L,R} = \frac{1}{2}(1 \pm \gamma_5) \psi \rightarrow \psi = \begin{pmatrix} u \\ d \\ s \end{pmatrix} \rightarrow \sqrt{\psi}^{SU(3)}$$

$$\mathcal{L} = \bar{\psi} i \not{D} \psi = \bar{\psi}_L i \not{D} \psi_L + \bar{\psi}_R i \not{D} \psi_R$$

$$\psi_L \rightarrow L \psi_L, \quad \psi_R \rightarrow R \psi_R \quad SU(3)_L \times SU(3)_R$$

$$\mathcal{L}_{QCD} = -\frac{1}{4} G^2 + \bar{\psi} i \not{D} \psi - \bar{\psi}_L \gamma_\mu \not{t}_a \psi_L - \bar{\psi}_R \gamma_\mu \not{t}_a \psi_R$$

$$- \bar{\psi}_L (S + iP) \psi_R - \bar{\psi}_R (S - iP) \psi_L$$

Real pure QCD $h_u = h_d = h_s = 0$, $S = M = \begin{pmatrix} m_u & & \\ & m_d & \\ & & m_s \end{pmatrix}$

+ QED

$$\rightarrow \begin{aligned} \mathcal{L}_m &= \mathcal{L}_m = -\frac{1}{2} \partial_\mu A_\nu \partial^\mu A^\nu \\ \mathcal{L}_m &= -\frac{1}{2} \partial_\mu \vec{W}_\nu \partial^\mu \vec{W}^\nu, \quad \mathcal{L}_m = 0 \end{aligned}$$

Invariance $\psi_L \rightarrow L(x) \psi_L, \quad \psi_R \rightarrow R(x) \psi_R$

$$\mathcal{L}_m \rightarrow L \mathcal{L}_m L^\dagger + (\partial_\mu L) L^\dagger$$

$$(S = \pm i p) \rightarrow L (S \pm i p) R^\dagger$$

~~2~~

$$\mathcal{L} = \frac{F^2}{4} \text{Tr} (D_\mu U D^\mu U^\dagger) + \frac{F^2 B_0}{2} \text{Tr} (S + i p) U^\dagger + (S - i p) U$$

$$D_\mu U = \partial_\mu U + i \overleftarrow{\mathcal{L}_m} U - i U \mathcal{R}_m$$

$$U = \exp i \frac{\lambda^a \phi^a}{F}$$

$\pi, K, \eta^8 \Rightarrow \vec{c} \cdot \vec{\pi}$

$$F \rightarrow \boxed{\pi \rightarrow W \rightarrow e \nu}$$

$$F = F_\pi = 93 \text{ MeV}$$

$$m_\pi^2 = B_0 (m_u + m_d)$$

$$m_K^2 = B_0 (m_s + m_u)$$

$$\frac{m_u + m_d}{m_s + m_u} = \frac{m_\pi^2}{m_K^2} \sim \frac{1}{25}$$

$$m_\pi = 133 \text{ MeV}, \quad m_K = 495 \text{ MeV}$$

$\mathcal{O}(E^4) \mathcal{L}$ - Gasser & Leutwyler 1984

$$L_i \sim \text{few} \times 10^{-3}$$

$$L_1 = 1.12 \pm 0.20 \times 10^{-3}$$

\vdots

l_1, l_2 from σ model
fits

$$\begin{aligned} \mathcal{L}_4 &= \sum_{i=1}^{10} L_i O_i \\ &= l_1 \left[\text{Tr} (D_\mu U D^\mu U^\dagger) \right]^2 + L_2 \text{Tr} (D_\mu U D_\nu U^\dagger) \cdot \text{Tr} (D^\mu U D^\nu U^\dagger) \\ &+ L_3 \text{Tr} (D_\mu U D^\mu U^\dagger D_\nu U D^\nu U^\dagger) \\ &+ L_4 \text{Tr} (D_\mu U D^\mu U^\dagger) \text{Tr} (\chi U^\dagger + U \chi^\dagger) \\ &+ L_5 \text{Tr} (D_\mu U D^\mu U^\dagger (\chi U^\dagger + U \chi^\dagger)) + L_6 \left[\text{Tr} (\chi U^\dagger + U \chi^\dagger) \right]^2 \\ &+ L_7 \left[\text{Tr} (\chi^\dagger U - U \chi^\dagger) \right]^2 + L_8 \text{Tr} (\chi U^\dagger \chi U^\dagger + U \chi^\dagger U \chi^\dagger) \\ &+ i L_9 \text{Tr} (L_{\mu\nu} D^\mu U D^\nu U^\dagger + R_{\mu\nu} D^\mu U^\dagger D^\nu U) + L_{10} \text{Tr} (L_{\mu\nu} U R^{\mu\nu} U^\dagger) \end{aligned}$$

$\sigma(\gamma\gamma \rightarrow \pi^0\pi^0)$

all loop

