

Variational Calculus

2nd law: $m\ddot{\eta}_i = -\vec{\nabla}V$ for a single body of mass m , in a conservative force $\vec{F} = -\vec{\nabla}V$

or in components $m\frac{d^2\eta_i}{dt^2} = -\frac{\partial V}{\partial \eta_i}$ (here: $i=x, y, z$)

now write: $m\ddot{\eta}_i = m\dot{v}_i = \frac{\partial}{\partial v_i} \left[\frac{1}{2} m \vec{v}^2 \right] = \frac{\partial}{\partial v_i} \left[\frac{1}{2} m \sum_{j=x,y,z} v_j^2 \right]$

$$\Rightarrow \frac{d}{dt} \left[\frac{\partial}{\partial \dot{\eta}_i} \left(\frac{1}{2} m \vec{v}^2 \right) \right] = -\frac{\partial V}{\partial \eta_i}$$

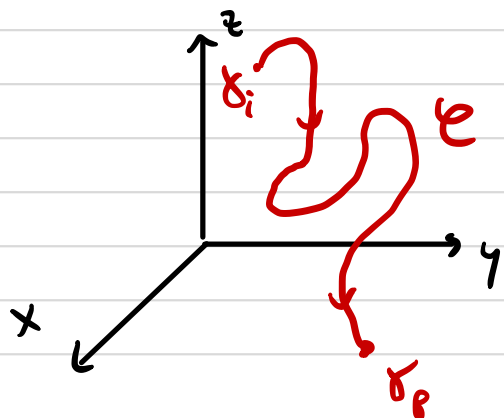
Let $L = T - V = \frac{1}{2} m \vec{v}^2 - V(\eta)$ ^{Lagrangian}

$$\text{Then: } \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\eta}_i} \right) - \frac{\partial L}{\partial \eta_i} = 0$$

$T = T(\dot{\eta}_i)$
 $V = V(\eta_i)$

We'll now see how this equation emerges in much more general contexts.

① Functionals and variational calculus



Curve: $(x(\gamma), y(\gamma), z(\gamma))$
parametrized by $\gamma \in [\gamma_i, \gamma_e]$

Functional:

Function $\beta(x) \mapsto \mathcal{F}[\beta] \in \mathbb{R}$
"Number"

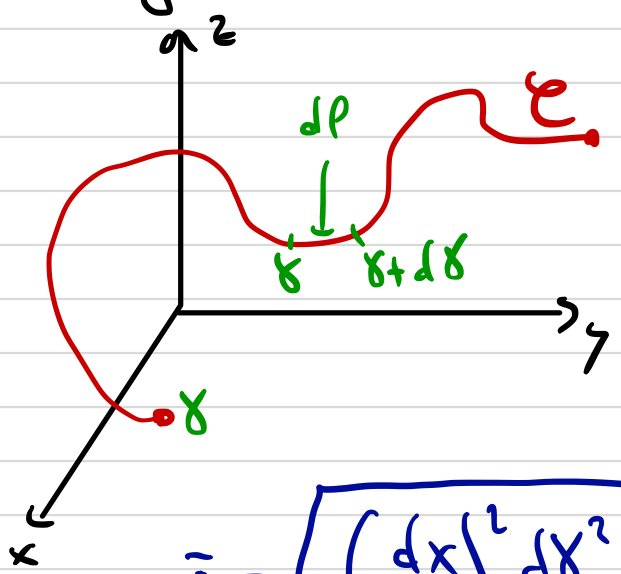
Ex: $\mathcal{F}[\beta] = \int_0^1 \beta(x) dx$ for every function defined on $[0, 1]$

$\mathcal{F}[\beta(x)=x] = \int_0^1 x dx = \left[\frac{x^2}{2}\right]_0^1 = \frac{1}{2}$

$\mathcal{F}[\beta(x)=x^3/2] = \frac{1}{2} \int_0^1 x^3 dx = \frac{1}{2} \left[\frac{x^4}{4}\right]_0^1 = \frac{1}{8}$

Ex: Curve $\mathcal{C}(\gamma) = (x_\gamma, y_\gamma, z_\gamma)$ w/ $x_\gamma = x(\gamma)$ etc.

Length of $\mathcal{C} =$ functional of $\mathcal{C}(\gamma)$



$dP = \|\vec{r}_{\gamma+d\gamma} - \vec{r}_\gamma\|$
 $= \sqrt{(x_{\gamma+d\gamma} - x_\gamma)^2 + (y_{\gamma+d\gamma} - y_\gamma)^2 + (z_{\gamma+d\gamma} - z_\gamma)^2}$

$= \sqrt{\left(\frac{dx}{d\gamma}\right)^2 d\gamma^2 + \left(\frac{dy}{d\gamma}\right)^2 d\gamma^2 + \left(\frac{dz}{d\gamma}\right)^2 d\gamma^2}$
 $= d\gamma \sqrt{\sum_{i=1}^d \left(\frac{dx_i}{d\gamma}\right)^2}$ $x_i = x, y, z$ for $d=3$
 $i=1, 2, 3$

Length:

$\mathcal{L} = \int_{\mathcal{C}} dP = \int_{\gamma_i}^{\gamma_f} d\gamma \sqrt{\sum_{i=1}^d \left(\frac{dx_i}{d\gamma}\right)^2}$

Variational Calculus:

Goal: find function(s) that maximize(s) (minimize(s)) a certain functional.

eg: Which curve gives the shortest path between two points?

$$I = \int_{x_i}^{x_f} F(\vec{x}(\gamma), \vec{x}'(\gamma)) d\gamma$$

$\vec{x} = (x_1, \dots, x_d)$
 $i=1, \dots, d.$

↑ Functional
↑ Function!

Just like to extremize a function $f(x)$, we find x such that $f'(x) = 0$ (derivative = 0), here we will extremize

$$I[\vec{x}(\gamma)] \text{ such that } \frac{\delta I}{\delta x_i(\gamma)} = 0 \quad i=1, \dots, d$$

↑ = Functional derivative
 $\equiv \sum_i \frac{\delta I}{\delta x_i} \delta x_i(\gamma)$

$$\vec{\tilde{x}}(\gamma) = \vec{x}(\gamma) + \delta \vec{x}(\gamma)$$

$$I[\vec{\tilde{x}}(\gamma)] = I[\vec{x}(\gamma)] + \int d\gamma \frac{\delta I}{\delta \vec{x}(\gamma)} \cdot \delta \vec{x}(\gamma) + \mathcal{O}(\delta \vec{x}^2)$$

↑ definition of $\frac{\delta I}{\delta x(\gamma)}$

"Taylor expansion" for functionals!

We also have:

$$\frac{\delta I}{\delta x_i(\gamma)} = \lim_{\delta x_i(\gamma) \rightarrow 0} \frac{I[\vec{x}(\gamma) + \delta \vec{x}(\gamma)] - I[\vec{x}(\gamma)]}{\delta x_i(\gamma)}$$

Let's compute this explicitly:

$$I[\bar{x}(\delta)] = \int_{\delta_i}^{\delta_b} F[\bar{x}(\delta) + \delta x(\delta), \bar{x}'(\delta) + \delta x'(\delta)] d\delta$$

Taylor expand

$$\approx \int_{\delta_i}^{\delta_b} \left(F[\bar{x}(\delta), \bar{x}'(\delta)] + \underbrace{\frac{\partial F}{\partial x_i}}_{\sum_{i=1}^d \frac{\partial F}{\partial x_i} \delta x_i} \delta x(\delta) + \frac{\partial F}{\partial x'_i} \delta x'(\delta) + \mathcal{O}(\delta x^2) \right) d\delta$$

$$= I[\bar{x}(\delta)] + \int_{\delta_i}^{\delta_b} d\delta \left[\frac{\partial F}{\partial x_i} \delta x(\delta) + \frac{\partial F}{\partial x'_i} \cdot \frac{d}{d\delta} \delta x(\delta) \right] + \dots$$



integrate by parts

boundary term = 0 if we assume $\delta x'(\delta_i) = \delta x'(\delta_b) = 0$

$$\Rightarrow I[\bar{x}'] - I[\bar{x}] = \int_{\delta_i}^{\delta_b} d\delta \left[\frac{\partial F}{\partial x_i} - \frac{d}{d\delta} \left(\frac{\partial F}{\partial x'_i} \right) \right] \delta x(\delta) + \dots$$

$$\Rightarrow \frac{\delta I}{\delta x(\delta)} = \frac{\partial F}{\partial x_i} - \frac{d}{d\delta} \left(\frac{\partial F}{\partial x'_i} \right) \quad \left(\bar{x}' = \frac{d}{d\delta} \bar{x} \text{ here} \right)$$

$I[\bar{x}(\delta)]$ is extremized by $\bar{x}(\delta)$ satisfying:

$$\frac{\partial F}{\partial x_i}(\bar{x}(\delta), \bar{x}'(\delta)) = \frac{d}{d\delta} \frac{\partial F}{\partial x'_i}(\bar{x}(\delta), \bar{x}'(\delta))$$

$i = 1, \dots, d$

Euler-Lagrange Equations

Two-line version (single variable):

$$\delta \left(\int d\gamma F(x, x') \right) = \int d\gamma \left(\frac{\partial F}{\partial x} \delta x + \frac{\partial F}{\partial x'} \delta x' \right)$$

↓
 $\frac{d}{d\gamma} \delta x$

$$= \int d\gamma \left(\frac{\partial F}{\partial x} - \frac{d}{d\gamma} \left(\frac{\partial F}{\partial x'} \right) \right) \delta x(\gamma)$$

$\frac{\delta I}{\delta x(\gamma)}$

Ⓐ An important theorem:

Let $E = \dot{x}' \cdot \frac{\partial F}{\partial \dot{x}'} - F$. Then
 with $\vec{x}(\gamma)$ such that $\delta I = 0$

$\frac{dE}{d\gamma} = - \frac{\partial F}{\partial \gamma}$

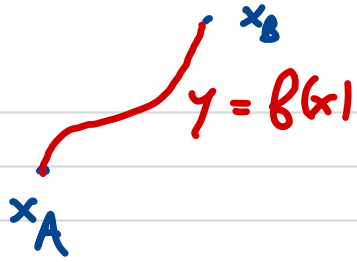
if $\frac{\partial F}{\partial \gamma} = 0$: $E = \text{Const}$ (first integral of the Euler-Lagrange equations)

Proof:

$$\begin{aligned} \frac{dE}{d\gamma} &= \frac{d}{d\gamma} \left(\dot{x}' \cdot \frac{\partial F}{\partial \dot{x}'} \right) - \frac{dF}{d\gamma} \\ &= \ddot{x}'' \cdot \frac{\partial F}{\partial \dot{x}'} + \dot{x}' \cdot \frac{d}{d\gamma} \left(\frac{\partial F}{\partial \dot{x}'} \right) - \left(\frac{\partial F}{\partial \gamma} + \frac{\partial F}{\partial \dot{x}'} \cdot \dot{x}'' + \frac{\partial F}{\partial x'} \cdot \dot{x}'' \right) \\ &= - \frac{\partial F}{\partial \gamma} \end{aligned}$$

$\frac{d}{d\gamma} \frac{\partial F}{\partial \dot{x}'}$

Example



Length of curve:

$$I = \int \sqrt{dx^2 + dy^2} = \int dx \sqrt{1 + (\beta')^2}$$

$$F(\beta, \beta') = \sqrt{1 + (\beta')^2}, \text{ indep of } \beta \text{ and } x$$

$$E = \beta' \frac{\partial F}{\partial \beta'} - F = \text{const} \Rightarrow \beta' = \text{const}$$

$$\Rightarrow \beta = Ax + B$$

A, B fixed by boundary conditions.

↑
shortest path
= straight line

$$\text{Euler Lagrange: } \frac{d}{dx} \left(\frac{\partial F}{\partial \beta'} \right) = \frac{\partial F}{\partial \beta} = 0 \Rightarrow \frac{\partial F}{\partial \beta'} = \text{const.}$$

Ⓜ Constraints (intro: we'll come back to this topic in the next chapter)

Example: (without functionals). Suppose we want to extremize (maximize in this case) $f(x, y) = x + y$ subject to the constraint $x^2 + y^2 = 1$.

Introduce: $F(x, y, \lambda) = x + y + \lambda(x^2 + y^2 - 1)$

↑ Lagrange multiplier

(7)

Extremize F : $\partial_x F = 1 + 2x\lambda = 0$

$$\partial_y F = 1 + 2y\lambda = 0$$

$$\partial_\lambda F = x^2 + y^2 - 1 = 0 : \text{constraint!}$$

$$\Rightarrow x = y = -\frac{1}{2\lambda} \text{ and } \frac{1}{2\lambda^2} = 1 \Rightarrow \lambda = \pm \frac{1}{\sqrt{2}}$$

$$\Rightarrow x = y = \pm \frac{\sqrt{2}}{2}, \quad \begin{array}{l} + \text{ sign: maximum} \\ - \text{ sign: minimum.} \end{array}$$

More generally: Functional $I[\vec{x}] = \int d\delta F(\vec{x}, \vec{x}')$

+ constraints: ① $Q_i[x] = \int d\delta q_i(\vec{x}(\delta)) = 0$
 $i = 1, \dots, \# \text{ constraints}$

② $q_i(\vec{x}(\delta)) = 0, \forall \delta.$

Modified functional: ① $\tilde{I} = I + \sum_i \lambda_i Q_i = I[x, \vec{\lambda}]$

② $\tilde{I} = \int d\delta \left(F + \sum_i \lambda_i(\delta) q_i(\vec{x}(\delta)) \right)$

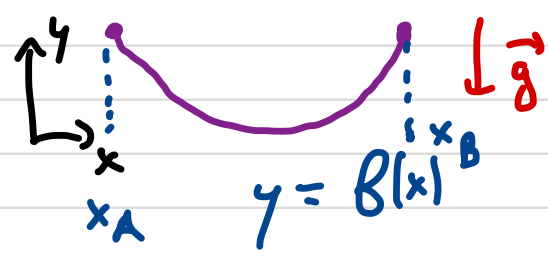
Euler Lagrange equations: $\hat{F} = F + \sum_i \lambda_i q_i$
Function of δ in case 2

$$\frac{\partial}{\partial x} \left(F + \sum_i \lambda_i q_i \right) = \frac{d}{d\delta} \left(\frac{\partial F}{\partial x'} \right) \Rightarrow \frac{d}{d\delta} \left(\frac{\partial F}{\partial x'} \right) - \frac{\partial F}{\partial x} = \sum_i \lambda_i q_i'$$

① $\frac{dI}{d\lambda_i} = 0 \Rightarrow Q_i = 0$: λ_i not a function here

② $\frac{\partial}{\partial \lambda_i} \tilde{F} = 0 \Rightarrow \frac{\partial \tilde{F}}{\partial \lambda_i(x)} = q_i(x) = 0$ (constraints)

Example



flexible cable suspended by two points.
(Catenary problem)

Minimize potential energy : $E = + g \overset{\text{mass/length}}{\downarrow} \int y \, dp$

forget units : minimize $I = \int_{x_A}^{x_B} dx \, \beta \sqrt{1 + (\beta')^2}$

with constraint : $Q = \int_{x_A}^{x_B} dx \sqrt{1 + (\beta')^2} - L$
(Length fixed!)

$$\tilde{I} = \int_{x_A}^{x_B} dx (\beta + \lambda) \sqrt{1 + (\beta')^2} - \lambda L$$

Let $\tilde{\beta} = \beta + \lambda$: extremize $\int \tilde{\beta} \sqrt{1 + (\tilde{\beta}')^2} dx$
 $F(\tilde{\beta}, \tilde{\beta}')$
indep. of x

you'll {extremize
minimize} an identical functional in the problem sets in the context of soap films.

See HW: 3 constants (including λ)

Fixed by boundary conditions $R = \beta(x_1)$
+ constraint $Q = 0$: 3 equations $= \beta(x_2)$