# Kepler's problem: Trajectory of a particle in a $1 / r$ potential 

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Our goal is to study the trajectory of a particle of mass $m$ in a central potential $U(r)=-k / r$ with $k>0$. Since the torque of the corresponding force is zero, the angular momentum is conserved, and the motion is effectively twodimensional (in the plane orthogonal to the initial angular momentum). We take the angular momentum to be along the $z$ axis, and use polar coordinates in the $x y$ plane. The conservation of angular momentum gives us that

$$
\begin{equation*}
\ell_{0}=m r^{2} \dot{\theta} \tag{1}
\end{equation*}
$$

is a conserved quantity. From the second law, we have the equation of motion

$$
\begin{equation*}
m\left(\ddot{r}-r \dot{\theta}^{2}\right)=-\frac{k}{r^{2}} \tag{2}
\end{equation*}
$$

Using the conservation law (1) to get rid of $\dot{\theta}$, we find the effective one-dimensional motion $m \ddot{r}=-d U_{\text {eff }} / d r$ with the effective potential:

$$
\begin{equation*}
U_{\mathrm{eff}}(r)=-\frac{k}{r}+\frac{\ell_{0}^{2}}{2 m r^{2}} \tag{3}
\end{equation*}
$$

The energy is given by

$$
\begin{equation*}
E=\frac{1}{2} m \dot{r}^{2}+U_{\mathrm{eff}}(r) \tag{4}
\end{equation*}
$$

We can also express the temporal derivative as

$$
\begin{equation*}
\dot{r}=\dot{\theta} \frac{d r}{d \theta}=\frac{\ell_{0}}{m r^{2}} \frac{d r}{d \theta} \tag{5}
\end{equation*}
$$

The simplest way to derive the shape of the trajectory is to use energy conservation and to introduce a new variable $u$ such that $r=1 / u$. We have

$$
\begin{equation*}
\dot{r}=\frac{\ell_{0} u^{2}}{m} \frac{d r}{d \theta}=\frac{\ell_{0} u^{2}}{m} \frac{d r}{d u} \frac{d u}{d \theta}=-\frac{\ell_{0}}{m} \frac{d u}{d \theta} \tag{6}
\end{equation*}
$$

Let's plug this expression into (4) (using $r=1 / u$ in (3))

$$
\begin{equation*}
E=\frac{\ell_{0}^{2}}{2 m}\left(\frac{d u}{d \theta}\right)^{2}-k u+\frac{\ell_{0}^{2}}{2 m} u^{2} \tag{7}
\end{equation*}
$$

Now, since energy is conserved along the trajectory, we have

$$
\begin{equation*}
\frac{d E}{d \theta}=0=\frac{d u}{d \theta}\left(\frac{\ell_{0}^{2}}{m} \frac{d^{2} u}{d \theta^{2}}-k+\frac{\ell_{0}^{2}}{m} u\right) \tag{8}
\end{equation*}
$$

which yields

$$
\begin{equation*}
\frac{d^{2} u}{d \theta^{2}}+u=\frac{k m}{\ell_{0}^{2}} \tag{9}
\end{equation*}
$$

This is the equation of a harmonic oscillator with a constant right hand side. The particular solution is that constant, so the general solution reads:

$$
\begin{equation*}
u=\frac{k m}{\ell_{0}^{2}}+A \cos \left(\theta-\theta_{0}\right) \tag{10}
\end{equation*}
$$

where $A$ and $\theta_{0}$ are integration constants. $\theta_{0}$ can be chosen by setting a reference angle from which angles are measured, so it is usually set to some convenient choice like $\theta_{0}=0$ or $\theta_{0}=\pi$. Going back to $r=1 / u$, we have

$$
\begin{equation*}
r(\theta)=\frac{\frac{\ell_{0}^{2}}{k m}}{1-\epsilon \cos \theta} \tag{11}
\end{equation*}
$$

where $\epsilon=-A \ell_{0}^{2} /(k m)$ and we have chosen $\theta_{0}=0$. This is the equation of a conic section in polar coordinates.

